

11/08/2017

# RAMAN EFFECT

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## Introduction :

If beam of monochromatic light passes through a transparent medium (solid, liquid, gases) then it gets scattered. The scattered radiation can be observed at right angles to the incident radiation. Generally the freq<sup>n</sup> of scattered radiation is same as that of incident radiation & is known as Rayleigh's scattering. But Smekel in 1923 & Raman and Heisenberg theoretically shows that scattered radiation should contain radiation at different freq<sup>n</sup> from the incident.

Raman & Krishan conducted a series of experimental work on scattering of light (visible radiation) by a various transparent substance & should find in 1928 that the scattered radiation do contain joint radiation with freq<sup>n</sup> greater as well as less than incident freq<sup>n</sup>.

This scattering effect is known as Raman effect. The spectral line in the scattered light with modified freq<sup>n</sup> are known as Raman lines.

When the photon of energy  $h\nu_i$  is incident on atom/molecule of scattering medium, the state of energy of atom/molecule of the medium changes from  $E_1$  to  $E_2$ . Hence the energy of scattered photon is given

by,

$$h\nu_s = h\nu_i - (E_2 - E_1) \quad h\nu_i = \text{energy of incident photon}$$

$$\nu_s = \nu_i - \frac{(E_2 - E_1)}{h}$$

$$\nu_s = \nu_i - \Delta\nu$$

where,  $\Delta\nu$  is Raman shift and

$$\Delta\nu = \nu_i - \nu_s$$



# RAMAN EFFECT

Wavelength	Frequency	Wavenumber	Energy	Momentum	Transition description
$\lambda$	$\nu$	$\tilde{\nu}$	$E = h\nu$	$p = \frac{h\nu}{c}$	Incident radiation
$\lambda_s$	$\nu_s$	$\tilde{\nu}_s$	$E_s = h\nu_s$	$p_s = \frac{h\nu_s}{c}$	Stokes line
$\lambda_a$	$\nu_a$	$\tilde{\nu}_a$	$E_a = h\nu_a$	$p_a = \frac{h\nu_a}{c}$	Anti-Stokes line

It should be noted that Raman shift ( $\Delta\nu$ ) does not depend upon the frequency of incident radiation but it depends on the nature of scatterer. Depending upon the value of  $\Delta\nu$ , Raman lines are classified as,

- (i)  $\Delta\nu > 0$  i.e. if  $\Delta\nu$  is +ve then the scattered lines are called Stokes line.
- (ii)  $\Delta\nu < 0$ , i.e. if  $\Delta\nu$  is -ve then the scattered lines are called anti-Stokes line.

As shown in fig:





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electronic field =  $F$   
dipole moment =  $\mu$  - magnetic moment

induced  
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### \* Characteristics of Raman line.

- 1) Raman lines are very weak in intensity as compared to Rayleigh's line.
- 2) Stokes line are more intense as compared to anti-stokes line.
- 3) Raman lines are symmetrically displaced about the undisplaced parent line (Rayleigh's scattered line).
- 4) Raman shift  $\Delta\nu$  generally lies in infra-red region of the spectrum.
- 5) The  $\Delta\nu$  represent the freq of infrared (vibrational freq) of the absorption band of the scattering material.
- 6) Even though  $\Delta\nu$  is in infrared region the Raman line ( $\nu_3$ ) appears in visible region.

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### \* Classical Theory of Raman Effect.

When atom or molecules are placed in external electric field  $F$ , the  $e^-$  & nuclei i.e.  $\pm ve$  charge centers gets displaced. Consequently, small electric dipoles are created, when medium is said to be polarised, as these dipoles align along the field direction.

The dipole moment ( $\mu$ ) of induced dipole is proportional to strength of electric field ( $F$ )

$$\mu \propto F$$

$$\mu = \alpha F$$

where proportionality const.  $\alpha$  is called polarisability of molecule.



When light is incident on transparent scattering medium, The mol. of medium are subjected to the electric field of electromagnetic radiation of freq<sup>n</sup>  $\nu$ ; which is given by,

$$F = F_0 \sin 2\pi\nu t \quad (3)$$

$$\therefore \mu = \alpha F = \alpha F_0 \sin 2\pi\nu t \quad (3)$$

Thus induced dipole moment oscillates with freq<sup>n</sup>  $\nu$ ; equal to freq<sup>n</sup> of incident radiation.

Let us now consider the vibrational & rotational interaction with incident radiation.

### Vibrational (Rotational) interaction ( $\nu_v$ )

If the molecule undergoes vibrational motion with freq<sup>n</sup>  $\nu_v$  then it introduces the periodical change in the polarisability of the mol. which is given as,

$$\alpha = \alpha_0 + \beta \sin 2\pi\nu_v t \quad (4)$$

where  $\alpha_0$  = equilibrium polarisability.

$\beta$  = rate of change of polarisability with

vibrations.

Now substituting  $\alpha$  from eq<sup>n</sup> (4) in eq<sup>n</sup> (3)

$$\mu = (\alpha_0 + \beta \sin 2\pi\nu_v t) F_0 \sin 2\pi\nu t \quad (5)$$

$$\mu = \alpha_0 F_0 \sin 2\pi\nu t + \beta F_0 \sin 2\pi\nu t \cdot \sin 2\pi\nu_v t$$

Using :  $\sin A \cdot \sin B = \frac{1}{2} (\cos(A-B) - \cos(A+B))$





$$\mu = \alpha_0 F_0 \sin 2\pi \nu_i t + \frac{1}{2} \beta F_0 [\cos 2\pi (\nu_i - \nu_v) t - \cos 2\pi (\nu_i + \nu_v) t] \quad (5)$$

eq<sup>n</sup> (5) shows that the induced dipole moment oscillate with freq<sup>n</sup>  $\nu_i$ ,  $(\nu_i - \nu_v)$  &  $(\nu_i + \nu_v)$

And hence freq<sup>n</sup> of emitted radiation called the scattered radiation are  $\nu_i$ ,  $(\nu_i - \nu_v)$  &  $(\nu_i + \nu_v)$ .

The scattered freq<sup>n</sup>  $\nu_s$  is equal to  $\nu_i$  correspond to Rayleigh's scattering, while the scattered freq<sup>n</sup>  $\nu_s$  correspond to Raman lines  $\nu_s = (\nu_i \pm \nu_v)$

Thus Raman shift  $\Delta \nu = \nu_v$  i.e. Raman shift correspond to freq<sup>n</sup> of IR-vibrational band. Raman shift represent the freq<sup>n</sup> of IR vibrational band.

2) Rotational interaction

When mole. is not optically isotropic. ∴ when the mole. rotate in the direction of electric field of incident radi. then polarisability of mole. varies periodically. In case of diatomic mole., the rotation through  $\pi^\circ$  brings the mole. in the same situation as in original. Hence the polarisability changes with freq<sup>n</sup> double that of rotational freq<sup>n</sup> ( $\nu_r$ ) of the mole. we can

express the polarisability  $\alpha$  as

$$\alpha = \alpha_0 + \beta' \sin (2\nu_r) 2\pi t$$

Now if put  $\alpha$  in eq<sup>n</sup> (7) in eq<sup>n</sup> (5).

$$\mu = (\alpha_0 + \beta' \sin 2\pi (2\nu_r) t) F_0 \sin 2\pi \nu_i t$$

$$\mu = \alpha_0 F_0 \sin 2\pi \nu_i t + \beta' F_0 \sin 2\pi (2\nu_r) t \cdot \sin 2\pi \nu_i t$$

$$\mu = \alpha_0 F_0 \sin 2\pi \nu_i t + \frac{1}{2} \beta' F_0 [\cos 2\pi (\nu_i - 2\nu_r) t - \cos 2\pi (\nu_i + 2\nu_r) t] \quad (8)$$



From eq<sup>n</sup> (8) it is clear that the Raman lines

I have  $\nu_{rot} = 20, 12, 20$   
 Hence  $\Delta\nu$  in the  $\nu_{rot}$  is given by

$$\Delta\nu = 2\nu_{rot}$$

Thus the Raman shift is the  $\nu_{rot}$  equal to twice the rotational  $\nu_{rot}$  of mole.

	Raman Spectra	IR - Spectra.
16/08/2017	1) It arises due to scattering of light by vibrating mole.	1) It arises due to absorption of light by vibrating mole.
	2) The observation of Raman Spectra depends upon polarising of the mole.	2) IR spectra are produced only if the mole possess permanent electric dipole moment.
	3) Raman lines weak in intensity & sample with conc. sol <sup>n</sup> are prepared	3) Generally dil. sol <sup>n</sup> are preferred.
	4) Water can be used as a solvent to prepare sol <sup>n</sup>	4) Water being opaque to IR it can't be used as solvent.
	5) optical system are made up of glass as Raman spectra observed in visible	5) optical system are made upto $CaF_2$ & $NaBr$ which are transparent to IR.

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6) Substance under investi-  
gating must be pure &  
colourless

6) These condition is not  
rigid.

7) vibrational freq<sup>s</sup> of large  
mole can be measured.

7) vibrational freq<sup>n</sup> of  
large mole. can't be  
measure

8) Even homonuclear diatomic  
mole. like  $H_2$ ,  $N_2$  etc  
produces Raman spectra

8) Homonuclear diatomic  
mole. can't produce  
IR spectra as did not  
possess permenent di-  
pole moment

9) Raman lines being very  
weak, the excitation sou-  
rce must have high in-  
tensity

9) It is not so